

Baxter's TQ-construction  
An alternative to the Bethe Ansatz

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## Schedule

- Baxter's TQ-Construction
- Application of TQ to the 6-vertex model
- (If there is time) Link to Bethe Ansatz results
- Extension to 8-vertex model

## Baxter's TQ-Construction for the 6-vertex model

- Given a Transfer Matrix  $\mathbf{T}$  belonging to Boltzmann weights  $a$ ,  $b$  and  $c$  we can find other transfer matrices  $\mathbf{T}'$  with corresponding Boltzmann weights  $a'$ ,  $b'$  and  $c'$  so that  $[\mathbf{T}, \mathbf{T}'] = 0$ . This will impose a (familiar) condition on the Boltzmann weights

$$\frac{a^2 + b^2 - c^2}{2ab} = \Delta = \Delta' = \frac{a'^2 + b'^2 - c'^2}{2a'b'}$$

- A fortiori, we can reparameterize  $a$ ,  $b$  and  $c$  in terms of a scaling factor  $\rho$  and rapidities  $\lambda$  and  $\phi$ . For a fixed choice of  $\Delta$  these transfer matrices form a 1-parameter family  $\mathcal{T} = \{\mathbf{T}(\lambda) | \lambda \in \mathbb{C}\}$ . All elements of a given  $\mathbf{T}(\lambda)$  are entire functions of  $\lambda$  (the spectral parameter).
- There exists a matrix function  $\mathbf{Q}(\lambda)$  that satisfies the commutation relations

$$[\mathbf{Q}(\lambda), \mathbf{Q}(\lambda')] = [\mathbf{Q}(\lambda), \mathbf{T}(\lambda')] = 0$$

and the TQ-relation

$$\mathbf{T}(\lambda)\mathbf{Q}(\lambda) = \sigma(\lambda - \phi)\mathbf{Q}(\lambda + 2\phi) + \sigma(\lambda + \phi)\mathbf{Q}(\lambda - 2\phi)$$

Where  $\sigma(\lambda) = [\rho \sinh \lambda]^L$

- The determinant of  $\mathbf{Q}(\lambda)$  is not entirely zero, and it decomposes into diagonal blocks on each subspace of fixed magnon number  $M$ .
- We recover the Bethe Ansatz and can now calculate the eigenvalues  $\Lambda(\lambda)$  of  $\mathbf{T}(\lambda)$ .

## Extension to 8-vertex model

- Given a Transfer Matrix  $\mathbf{T}$  belonging to Boltzmann weights  $a, b, c$  and  $d$  we can find other transfer matrices  $\mathbf{T}'$  with corresponding Boltzmann weights  $a', b', c'$  and  $d'$  so that  $[\mathbf{T}, \mathbf{T}'] = 0$ . This will impose a conditions on the Boltzmann weights

$$\frac{a^2 + b^2 - c^2 - d^2}{2(ab + cd)} = \Delta = \Delta' = \frac{a'^2 + b'^2 - c'^2 - d'^2}{2(a'b' + c'd')}$$

$$\frac{ab - cd}{ab + cd} = \Gamma = \Gamma' = \frac{a'b' - c'd'}{a'b' + c'd'}$$

- A fortiori, we can reparameterize  $a, b$  and  $c$  in terms of a scaling factor  $\rho$  and **three other numbers**  $\lambda, \kappa$  and  $\phi$ . For a fixed choice of  $\Delta$  and  $\Gamma$  these transfer matrices form a 1-parameter family  $\mathcal{T} = \{\mathbf{T}(\lambda) | \lambda \in \mathbb{C}\}$ . All elements of a given  $\mathbf{T}(\lambda)$  are entire functions of  $\lambda$  (the spectral parameter).
- There exists a matrix function  $\mathbf{Q}(\lambda)$  that satisfies the commutation relations

$$[\mathbf{Q}(\lambda), \mathbf{Q}(\lambda')] = [\mathbf{Q}(\lambda), \mathbf{T}(\lambda')] = 0$$

and the TQ-relation

$$\mathbf{T}(\lambda)\mathbf{Q}(\lambda) = \sigma(\lambda - \phi)\mathbf{Q}(\lambda + 2\phi) + \sigma(\lambda + \phi)\mathbf{Q}(\lambda - 2\phi)$$

Where  $\sigma(\lambda)$  is a significantly more complicated function.

- The determinant of  $\mathbf{Q}(\lambda)$  is not entirely zero, and it decomposes into diagonal blocks on **the subspace of even and the subspace of odd magnon numbers**.
- We recover equations **similar to the Bethe Ansatz** and can now calculate the eigenvalues  $\Lambda(\lambda)$  of  $\mathbf{T}(\lambda)$ .